## 轨道空间上的Frobenius流形结构

左达峰
Joint with B．Dubrovin and Youjin Zhang

School of mathematics<br>Korea Institute for Advanced Study

中国科技大学数学系

Chern Institute of Mathematics
Nov． 222007

## Main References

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3. Dafeng Zuo, International Mathematics Research Notices 8(2007)rnm020-25
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Extended affine Weyl groups and Frobenius manifolds-II
The first draft: math.DG/0502365
5.-, Geometric structures related to new extended affine Weyl groups, in preparation.

## Outline of this talk

§1. Physical background [Ref.1]
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§1. Physical background

## 2-dimensional Toplogical field theory (2D TFT)

QFT on $\Sigma$ consists of

- Local fields $\phi_{\alpha}(x), x \in \Sigma$, eg. gravity field: the metric $g_{j j}(x)$
- Classical action

$$
S[\phi]=\int_{\Sigma} L\left(\phi, \phi_{x}, \cdots\right)
$$

Remark. Classical field theory is determined by the Euler-Lagrangian equations $\frac{\delta S}{\delta \phi_{\alpha}(x)}=0$.

- Partition function

$$
Z_{\Sigma}=\int[d \phi] e^{-S[\phi]}
$$

- Correlators

$$
\left\langle\phi_{\alpha}(x) \phi_{\beta}(y) \cdots\right\rangle=\int[d \phi] \phi_{\alpha}(x) \phi_{\beta}(y) \cdots e^{-S[\phi]}
$$

Topological invariance

$$
\frac{\delta S}{\delta g_{i j}(x)}=0 \quad\left(i . e ., \delta g_{i j}(x)=\text { arbitrary, } \delta S=0\right)
$$

Remark. conformal field theory: $\delta g_{i j}(x)=\epsilon g_{i j}(x), \delta S=0$.
$\Rightarrow$ correlators are numbers depending only on the genus $g=g(\Sigma)$

$$
\left\langle\phi_{\alpha}(x) \phi_{\beta}(y) \cdots\right\rangle=\left\langle\phi_{\alpha} \phi_{\beta} \cdots\right\rangle_{g}
$$

Example. 2D gravity with Hilbert-Einstein action

$$
S=\frac{1}{2 \pi} \int R \sqrt{g} d^{2} x=\chi(\Sigma)
$$

There are two ways of quantization of this functional to obtain 2D quantum gravity.

1. [Matrix gravity] Base on an approximate discrete version of the model ( $\Sigma \rightarrow$ Polyhedron) $\rightsquigarrow$ Matrix integrals

$$
Z_{N}(t)=\int_{X=X^{*}} e^{-\operatorname{tr}\left(X^{2}+t_{1} X^{4}+t_{2} X^{6}+\cdots\right)} d X
$$

$N \rightarrow \infty \rightsquigarrow \tau$-function of KdV hierarchy $\rightsquigarrow$ a solution of 2D gravity
2.[Topological 2D gravity] Base on an approximate supersymmetric extension of Hilbert-Einstein Lagrangian $\rightsquigarrow$

$$
\sigma_{p} \leftrightarrow c_{p} \in H_{*}\left(\mathcal{M}_{g, n}\right)
$$

and the genus g correlators of the topological gravity are expressed as

$$
\left\langle\sigma_{p_{1}} \cdots \sigma_{p_{n}}\right\rangle=\#\left(c_{p_{1}} \cap \cdots c_{p_{n}}\right)=\prod_{i=1}^{n}\left(2 p_{i}+1\right)!!\int_{\overline{\mathcal{M}}_{g, n}} \psi_{1}^{p_{1}} \cdots \psi_{n}^{p_{n}},
$$

where $\psi_{i}=c_{1}\left(L_{i}\right) \in H^{*}\left(\overline{\mathcal{M}}_{g, n}\right)$ (the first Chern classes).

Witten conjecture [Proved by Kontsevich]
$\int_{\overline{\mathcal{M}}_{g, n}} \psi_{1}^{p_{1}} \cdots \psi_{n}^{p_{n}}$ is governed by the $\tau$-function of KdV hierarchy.

Problem. To find a rigorous mathematical foundation of 2D topological field theory.
M.F.Atiyah, Publ.Math. IHES. 68(1988)175-186. (inspired by G.Segal's axiomatization of CFT), for arbitrary dimension

Matter sector of a 2D TFT is specified by

1. $\mathcal{A}=$ the space of local physical states, $\operatorname{dim} \mathcal{A}<\infty$ basis $\left\{\phi_{1}=1, \cdots, \phi_{n}\right\}$ (primary observables)
2. an assignment

$$
(\Sigma, \partial \Sigma) \mapsto v_{(\Sigma, \partial \Sigma)} \in A_{(\Sigma, \partial \Sigma)}
$$

for any oriented 2-surface $\Sigma$ with an oriented boundary $\partial \Sigma$ depends only on the topology of the the pair $(\Sigma, \partial \Sigma)$

$$
\begin{aligned}
& A_{(\Sigma, \partial \Sigma)}:=\left\{\begin{array}{cl}
\mathbb{C}, & \text { if } \quad \partial \Sigma=\emptyset ; \\
A_{1} \otimes \cdots \otimes A_{k}, & \text { if } \quad \partial \Sigma=C_{1} \cup \cdots \cup C_{k}
\end{array}\right. \\
& A_{i}:= \begin{cases}\mathcal{A}, & \text { oriention of } C_{i} \text { is coherent to that of } \Sigma ; \\
\mathcal{A}^{*}, & \text { otherwise }\end{cases}
\end{aligned}
$$

The assignment satisfies three axioms: (see the attached files)

1. Normalization; 2. Multiplicativity; 3. Factorization.

Denote a symmetric polylinear function on the space of the states by

$$
v_{g, s}:=v_{(\Sigma, \partial \Sigma)} \in \underbrace{\mathcal{A}^{*} \otimes \cdots \otimes \mathcal{A}^{*}}_{s}, \quad g=g(\Sigma)
$$

The genus g correlators of the fields $\phi_{\alpha_{1}}, \cdots, \phi_{\alpha_{s}}$ are defined by

$$
\left\langle\phi_{\alpha_{1}} \cdots \phi_{\alpha_{s}}\right\rangle_{g}:=v_{g, s}\left(\phi_{\alpha_{1}} \otimes \cdots \otimes \phi_{\alpha_{s}}\right)
$$

Theorem[Dijgraff etc.] $\mathcal{A}$ carries a natural structure of a Frobenius algebra $\left(\mathcal{A}, \bullet,\langle\rangle,, \phi_{1}\right)$. All correlators can be expressed in a pure algebraic way in terms of this algebra, i.e.,

$$
\left\langle\phi_{\alpha_{1}} \cdots \phi_{\alpha_{s}}\right\rangle_{g}=\left\langle\phi_{\alpha_{1}} \bullet \ldots \bullet \phi_{\alpha_{s}}, H^{g}\right\rangle
$$

where $H=\eta^{\alpha \beta} \phi_{\alpha} \bullet \phi_{\beta}$ and $\eta_{\alpha \beta}=\left\langle\phi_{\alpha}, \phi_{\beta}\right\rangle$.

Definition. A Frobenius algebra is a pair $(\mathcal{A}, \bullet,\langle\rangle, e$,$) satisfying$

1. $\mathcal{A}$ is a commutative and associative algebra over $\mathcal{K}$ with a unit $e$;
2. $\langle\rangle:, \mathcal{A} \times \mathcal{A} \rightarrow \mathcal{K}$ is a symmetric non-degenerate inner product;
3. $\langle$,$\rangle is invariant, i.e., \langle a \bullet b, c\rangle=\langle a, b \bullet c\rangle$.

Example A[Topological sigma model]. X a smooth projective variety,

$$
\operatorname{dim}_{\mathbb{C}} X=d, H^{\text {odd }}(X)=0, \mathcal{A}=H^{*}(X), \operatorname{dim} \mathcal{A}=n
$$

primary observables $\leftrightarrow$ cohomologies $\phi_{1}=1, \cdots, \phi_{n} \in H^{*}(X)$

$$
\left\langle\phi_{i}, \phi_{j}\right\rangle:=\int_{X} \phi_{i} \cup \phi_{j}
$$

Claim: $\left(\mathcal{A}, \cup,\langle\rangle,, \phi_{1}\right)$ is a Frobenius algebra.

Example B[Topological Landau-Ginsburg model]. Let $f(x)$ be an analytic function, $x=\left(x_{1}, \cdots, x_{N}\right) \in \mathbb{C}^{N}$ with an isolated singularity at $x=0$ of the multiplicity $n$, i.e., $\left.d f\right|_{x=0}=0$.

$$
\begin{array}{r}
\mathcal{A}:=\mathbb{C}\left[x_{1}, \cdots, x_{N}\right] /\left(\frac{\partial f}{\partial x_{1}}, \cdots, \frac{\partial f}{\partial x_{N}}\right), \quad \operatorname{dim} \mathcal{A}=n \\
\text { primary observables } \leftrightarrow \phi_{1}=1, \phi_{2}(x), \cdots, \phi_{n}(x) \in \mathcal{A} \\
\langle\phi(x), \psi(x)\rangle:=\frac{1}{(2 \pi i)^{N}} \int_{\cap\left|\frac{\partial f}{\partial x_{i}}\right|=\epsilon} \frac{\phi(x) \psi(x)}{\frac{\partial f}{\partial x_{1}} \cdots \frac{\partial f}{\partial x_{N}}} d^{N} x
\end{array}
$$

Claim: $\left(\mathcal{A},\langle\rangle,, \phi_{1}, \cdot\right)$ is a Frobenius algebra.

Next, they consider a particular topological perturbation, which preserves the topological invariance:

$$
\begin{gathered}
S \mapsto \tilde{S}(t):=S-\sum_{\alpha=1}^{n} t^{\alpha} \int_{\Sigma} \Omega \\
\left\langle\phi_{\alpha}(x) \phi_{\beta}(y) \cdots\right\rangle(t) \equiv \int[d \phi] \phi_{\alpha}(x) \phi_{\beta}(y) \cdots e^{-\tilde{S}(t)}
\end{gathered}
$$

Theorem. [WDVV, 1991] The perturbed Frobenius algebra $\mathcal{A}_{t}$ satisfies WDVV equations of associativity

$$
\frac{\partial^{3} F}{\partial t^{\alpha} \partial t^{\beta} \partial t^{\lambda}} \eta^{\lambda \mu} \frac{\partial^{3} F}{\partial t^{\mu} \partial t^{\delta} \partial t^{\gamma}}=\frac{\partial^{3} F}{\partial t^{\delta} \partial t^{\beta} \partial t^{\lambda}} \eta^{\lambda \mu} \frac{\partial^{3} F}{\partial t^{\mu} \partial t^{\alpha} \partial t^{\gamma}},
$$

with a quasihomogeneity condition

$$
\mathcal{L}_{E} F=(3-d) F+\text { quadratic polynomial in } t
$$

where

$$
E=t^{1} \partial_{1}+\text { linear in } t^{2}, \cdots, t^{n}
$$

is Euler vector field and $\phi_{1}$ is unit,

$$
\eta_{\alpha \beta}=\left\langle\phi_{\alpha} \phi_{\beta}\right\rangle_{0}(t)=\text { const., } \quad\left\langle\phi_{\alpha} \phi_{\beta} \phi_{\gamma}\right\rangle_{0}(t)=\frac{\partial^{3} F}{\partial t^{\alpha} \partial t^{\beta} \partial t^{\gamma}}
$$

for some function $F(t)$, called primary free energy.
[B.Dubrovin's idea,1992]: Add the above statement [WDVV] as a new axiom of 2D TFT. That is to say, to reconstruct the building of 2D TFT on the base of WDVV equations.

Example $\mathbf{A}^{\prime}$. Frobenius algebra $\mathcal{A}_{t}$ : quantum cohomology of $X$

$$
t=\left(t^{\prime}, t^{\prime \prime}\right), t^{\prime} \in H^{2}(X) / 2 \pi i H^{2}(X, \mathbb{Z}), t^{\prime \prime} \in H^{* \neq 2}(X)
$$

The primary energy $F(t)$ is the generating function of the genus 0 Gromov-Witten invariants.

Particularly, $X=\mathbb{C} \mathbb{P}^{1}$,
Quantum cohomology of $\mathbb{C P}^{1}=\mathbb{C}[\phi] /\left(\phi_{2}=e^{t^{2}}\right)$

$$
F(t)=\frac{1}{2}\left(t^{1}\right)^{2}\left(t^{2}\right)+e^{t^{2}}, E=t^{1} \partial_{1}+2 \partial_{2}, e=\partial_{1}
$$

Example B'. Set $f_{t}(x)=f(x)+\sum_{\alpha=1}^{n} s^{\alpha}(t) \phi_{\alpha}(x)$, then the deformed Frobenius algebra $\mathcal{A}_{t}$ is

$$
\mathcal{A}_{t}:=\mathbb{C}\left[x_{1}, \cdots, x_{N}\right] /\left(\frac{\partial f_{t}}{\partial x_{1}}, \cdots, \frac{\partial f_{t}}{\partial x_{N}}\right)
$$

primary observables $\leftrightarrow$ elements of the Jacobi ring $\mathcal{A}$

$$
\langle\phi(x), \psi(x)\rangle \equiv \eta_{i j}(t):=\frac{1}{(2 \pi i)^{N}} \int_{\cap\left|\frac{\partial f_{t}}{\partial x_{i}}\right|=\epsilon} \frac{\phi(x) \psi(x)}{\frac{\partial f_{t}}{\partial x_{1}} \cdots \frac{\partial f_{t}}{\partial x_{N}}} d^{N} x
$$

Particularly, $f(x)=x^{n+1}$
The simple singularity of type $A_{n}, \mathcal{A}=\mathbb{C}[x] /\left(x^{n}\right)$.
§2. Definitions and Examples

Back to the main problem: we have a family of Frobenius algebras $\mathcal{A}_{t}$ depending on the parameters $t=\left(t^{1}, \cdots, t^{n}\right)$. Write

$$
M=\text { the space of parameters }
$$

and we have a fiber bundle

$$
\stackrel{\downarrow \mathcal{A}_{t}}{t \in M^{\prime}}
$$

The basic idea is to identify this fiber bundle with the tangent bundle TM of the manifold $M$.

Definition. A Frobenius structure of charge $d$ on $M$ is the data $(M, \bullet,\langle\rangle, e, E$,$) satisfying$
(i) $\eta:=\langle$,$\rangle is a flat pseudo-Riemannian metric and \nabla e=0$;
(ii) $\left(T_{m} M, \bullet, \eta, e\right)$ is a Frobenius algebra which depends smoothly on $m \in M$;
(iii) $\left(\nabla_{w} c\right)(x, y, z)$ is symmetric, where $c(x, y, z):=\langle x \bullet y, z\rangle$;
(iv) A linear vector field $E \in \operatorname{Vect}(M)$ must be fixed on $M$, i.e. $\nabla \nabla E=0$ such that

$$
\mathcal{L}_{E}\langle,\rangle=(2-d)\langle,\rangle, \quad \mathcal{L}_{E} \bullet=\bullet, \quad \mathcal{L}_{E} e=-e .
$$

Theorem. [B.Dubrovin 1992] There is a one to one correspondence between a Frobenius manifold and the solution $F(\mathbf{t})$ of WDVV equations of associativity

$$
\frac{\partial^{3} F}{\partial t^{\alpha} \partial t^{\beta} \partial t^{\lambda}} \eta^{\lambda \mu} \frac{\partial^{3} F}{\partial t^{\mu} \partial t^{\delta} \partial t^{\gamma}}=\frac{\partial^{3} F}{\partial t^{\delta} \partial t^{\beta} \partial t^{\lambda}} \eta^{\lambda \mu} \frac{\partial^{3} F}{\partial t^{\mu} \partial t^{\alpha} \partial t^{\gamma}},
$$

with a quasihomogeneity condition

$$
\mathcal{L}_{E} F=(3-d) F+\text { quadratic polynomial in } \mathrm{t} .
$$

Definition. A Frobenius manifold is called semisimple if the algebra $\left(T_{m} M, \bullet\right)$ are semisimple at generic $m$.

Main mathematical applications of Frobenius manifolds
$\star$ The theory of Gromov - Witten invariants,
$\star$ Singularity theory,
$\star$ Hamiltonian theory of integrable hierarchies,
$\star$ Differential geometry of the orbit spaces of reflection groups and of their extensions $\rightsquigarrow$ semisimple Frobenius manifolds.

Definition. An intersection form of Frobenius manifold is a symmetric bilinear form on the cotangent bundle $T^{*} M$ defined by

$$
\left(\omega_{1}, \omega_{2}\right)^{*}=i_{E}\left(\omega_{1} \cdot \omega_{2}\right), \quad \omega_{1}, \omega_{2} \in T^{*} M
$$

Here the multiplication law on the cotangent planes is defined using the isomorphism

$$
\langle,\rangle: T M \rightarrow T^{*} M
$$

The discriminant $\Sigma$ is defined by

$$
\Sigma=\left\{t|\operatorname{det}(,)|_{T_{t}^{*} M}=0\right\} \subset M .
$$

Theorem. [B.Dubrovin 1992]
The metrics $\eta:=\langle$,$\rangle and g:=(,)^{*}$ form a flat pencil on $M \backslash \Sigma$, i.e.,

1. The metric $h^{\alpha \beta}=\eta^{\alpha \beta}+\lambda g^{\alpha \beta}$ is flat for arbitrary $\lambda$ and
2. The Levi-Civita connection for the metric $h^{\alpha \beta}$ has the form

$$
\Gamma_{\delta_{(h)}}^{\alpha \beta}=\Gamma_{k_{(\eta)}}^{\alpha \beta}+\lambda \Gamma_{k_{(g)}}^{\alpha \beta},
$$

where $\Gamma_{\delta_{(h)}}^{\alpha \beta}=-h^{\alpha \gamma} \Gamma_{\delta \gamma_{(h)}}^{\beta}, \Gamma_{\delta_{(g)}}^{\alpha \beta}=-g^{\alpha \gamma} \Gamma_{\delta \gamma_{(g)}}^{\beta}, \Gamma_{\delta_{(\eta)}}^{\alpha \beta}=-\eta^{\alpha \gamma} \Gamma_{\delta \gamma_{(\eta)}}^{\beta}$.

The holonomy of the local Euclidean structure defined on $M \backslash \Sigma$ by the intersection form $(,)^{*}$ gives a representation

$$
\mu: \pi_{1}(M \backslash \Sigma) \rightarrow \operatorname{Isometries}\left(\mathbb{C}^{n}\right)
$$

Definition. The group

$$
W(M):=\mu\left(\pi_{1}(M \backslash \Sigma)\right) \subset \operatorname{Isometries}\left(\mathbb{C}^{n}\right)
$$

is called a monodromy group of Frobenius manifold.

$$
\Longrightarrow \quad M \backslash \Sigma=\Omega / W(M), \quad \Omega \subset \mathbb{C}^{n}
$$

[B.Dubrovin's conjecture] The monodromy group is a discrete group for a solution of WDVV equations with good properties.

Example. $\left[W(M)=\right.$ Coxeter group $\left.A_{1}\right] n=1, M=\mathbb{R}, t=t^{1}$,

$$
F(t)=\frac{1}{6} t^{3}, \quad E=t \partial_{t}, \quad e=\partial_{t}, \quad \eta^{11}=<\partial_{t}, \partial_{t}>=1
$$

$\rightsquigarrow$ dispersionless KdV hierarchy $\rightsquigarrow$ Witten Conjecture.

Example. $\left[W(M)=\right.$ extended affine Weyl group $\left.\widetilde{W}\left(A_{1}\right)\right]$
Quantum cohomology of $\mathbb{C P}^{1}$ :

$$
F=\frac{1}{2}\left(t^{1}\right)^{2} t^{2}+e^{t^{2}}, E=t^{1} \partial_{1}+2 \partial_{2}, e=\partial_{1}
$$

$\rightsquigarrow$ dispersionless extended Toda hierarchy $\rightsquigarrow$ Toda Conjecture.

Question 1. Given a Frobenius manifold, how to find the monodromy group? (Some cases can be computed).
Question 2. Which kind of groups can be served as the monodromy groups of some Frobenius manifolds?
\& Coxeter groups [B.Dubrovin1996]
\& Extended affine Weyl groups [B.Dubrovin Youjin Zhang 1996]
[Dubrovin-Zhang-Zuo 2005,general], [2007,new cases]
For the general case of type $E$, still open?
\& Jacobi forms $J\left(A_{n}\right), J\left(B_{n}\right), J\left(G_{2}\right)[\mathrm{n}=1$, B.Dubrovin 1996, general $n$, M.Bertola 2000], $J\left(E_{6}\right), J\left(D_{4}\right)$ [Satake.I 1993, 1998] Open for the rest?
\& Elliptic Weyl groups [Satake.I 2006, math.AG/0611553]
§3 Frobenius manifolds and Coxeter groups

Let W be a finite irreducible Coxeter group.

$$
W \curvearrowright V \quad W \curvearrowright S(V)
$$

[Chevalley Theorem]. The ring $S(V)^{W}$ of $W$-invariant polynomial functions on $V$

$$
\mathbb{C}\left[x_{1}, \cdots, x_{n}\right]^{W} \simeq \mathbb{C}\left[y^{1}, \cdots, y^{n}\right]
$$

where $y^{i}=y^{i}\left(x_{1}, \cdots, x_{n}\right)$ are certain homogeneous $W$-invariant polynomials of degree $\operatorname{deg} y^{i}=d_{i}, i=1, \cdots, n$.

The maximal degree $h$ is called the Coxeter number. We use the ordering of the invariant polynomials

$$
\operatorname{deg} y^{n}=d_{n}=h>d_{n-1}>\cdots>d_{1}=2
$$

The degrees satisfy the duality condition

$$
d_{i}+d_{n-i+1}=h+2, \quad i=1, \cdots, n .
$$

$$
W \curvearrowright V \quad W \curvearrowright V \otimes \mathbb{C}
$$

## $\mathcal{M}=V \otimes \mathbb{C} / W \quad$ affine algebraic variety

$$
S(V)^{W} \quad \text { the coordinate ring of } \mathcal{M}
$$

$V \rightsquigarrow$ flat manifold $\quad\left(V,\left\{x_{1}, \cdots, x_{n}\right\},\left(d x_{a}, d x_{b}\right)^{*}=\delta_{a b}\right)$
$\rightsquigarrow\left(\mathcal{M} \backslash \Sigma, g^{i j}(y)\right)$

$$
g^{i j}(y):=\left(d y^{i}, d y^{j}\right)^{*}=\sum_{a, b=1}^{n} \frac{\partial y^{i}}{\partial x_{a}} \frac{\partial y^{j}}{\partial x_{b}} \delta_{a b}
$$

Lemma.[K.Saito etc 1980]

1. The metric $\left(g^{i j}(y)\right)$ is flat on $\mathcal{M} \backslash \Sigma$.
2. These $g^{i j}(y)$ are at most linear w.r.t $y^{n}$.

Write

$$
e:=\frac{\partial}{\partial y^{n}} .
$$

Introduce a new metric,

$$
\eta^{i j}(y):=\left\langle d y^{i}, d y^{j}\right\rangle=\mathcal{L}_{e} g^{i j}(y)=\frac{\partial g^{i j}(y)}{\partial y^{n}} .
$$

Theorem. [K.Saito etc. 1980, B.Dubrovin 1992]
The metrics $\langle$,$\rangle and (, )* form a flat pencil of metrics.$
Moreover, there exist homogeneous polynomials

$$
t^{1}(x), \cdots, t^{n}(x)
$$

of degrees $d_{1}, \cdots, d_{n}$ respectively such that the matrix

$$
\left\langle d t^{i}, d t^{j}\right\rangle:=\eta^{i j}=\frac{\partial g^{i j}(t)}{\partial t^{n}}
$$

is a constant nondegenerate matrix.

Theorem.[B.Dubrovin, 1992] There exists a unique Frobenius structure of charge $d=1-\frac{2}{h}$ on the orbit space $\mathcal{M}$ polynomial in $t^{1}, t^{2}, \cdots, t^{n}$ such that

1. The unity vector field e coincides with $\frac{\partial}{\partial y^{n}}=\frac{\partial}{\partial t^{n}}$;
2. The Euler vector field has the form

$$
E=\sum_{\alpha=1}^{n} d_{\alpha} t^{\alpha} \frac{\partial}{\partial t^{\alpha}}
$$

Theorem. [B.Dubrovin's conjecture, 1996. C.Hertling, 1999]
Any irreducible semisimple polynomial Frobenius manifold with positive invariant degrees is isomorphic to the orbit space of a finite Coxeter group.

Our question and result I
Lemma. [M.Bertola, 1998] For $B_{n}$ and $1 \leq k \leq n$,

1. These $g^{i j}(y)$ are at most linear w.r.t $y^{k}$
2. The space $\mathcal{M}$ carries a flat pencil of metrics

$$
\begin{equation*}
g^{i j}(y) \text { and } \eta^{i j}(y)=\frac{\partial g^{i j}(y)}{\partial y^{k}} \tag{0.1}
\end{equation*}
$$

Question: If $k \neq n$, how to construct the flat coordinates of $\eta^{i j}(y)$ and the corresponding Frobenius manifolds?
M.Bertola's results (unpublished 1999)
M.Bertola started from the superpotential

$$
\lambda(p)=p^{-2(n-k)}\left(\sum_{a=1}^{n} p^{2(n-a)} y_{a}+p^{2 n}\right)
$$

to compute the corresponding potential $F(t)$ and obtained $n$ different Frobenius structures related to $B_{n}$. For example,

$$
\eta\left(\partial^{\prime}, \partial^{\prime \prime}\right)=-\sum_{|\lambda|<\infty} \operatorname{res}_{d \lambda=0} \frac{\partial^{\prime}(\lambda(p) d p) \partial^{\prime \prime}(\lambda(p) d p)}{d \lambda(p)}
$$

## Our construction is different.

We started from the flat pencil of metrics.
The first step is to construct the flat coordinate $t^{1}, \cdots, t^{n}$.
The second step is to show that $g^{i j}(t)$ and the $\Gamma_{m}^{i j}(t)$ are weighted homogeneous polynomials of $t^{1}, \ldots, t^{n}, \frac{1}{t^{n}}$.
The last step is to get the Frobenius structure.

Write

$$
\tilde{d}_{j}=\frac{j}{k}, \quad j \leq k, \quad \tilde{d}_{m}=\frac{2 k(n-m)+1}{2 k(n-k)}, \quad m>k
$$

Main Theorem.[Zuo IMRN-2007] For any fixed integer
$1 \leq k \leq n$, there exists a unique Frobenius structure of charge
$d=1-\frac{1}{k}$ on the orbit space $\mathcal{M} \backslash\left\{t^{n}=0\right\}$ of $B_{n}$ (or $D_{n}$ )
polynomial in $t^{1}, t^{2}, \cdots, t^{n}, \frac{1}{t^{n}}$ such that

1. The unity vector field e coincides with $\frac{\partial}{\partial y^{k}}=\frac{\partial}{\partial t^{k}}$;
2. The Euler vector field has the form

$$
E=\sum_{\alpha=1}^{n} \tilde{d}_{\alpha} t^{\alpha} \frac{\partial}{\partial t^{\alpha}}
$$

Theorem. [Zuo IMRN-2007] There is an isomorphism between them.

Motived by James T.Ferguson and I.A.B. Strachan' work, Logarithmic deformations of the rational superpotential/Landau-Ginzburg constructions of solutions of the WDVV equations, arXiv:Math-ph/0605078 we consider a water-bag reduction as follows

$$
\lambda(p)=p^{-2(n-k)}\left(\sum_{a=1}^{n} p^{2(n-a)} y_{a}+p^{2 n}\right)+\sum_{i=1}^{M} k_{i} \log \left(p^{2}-b_{i}^{2}\right) .
$$

Remark. Don't determine a full Frobenius manifold because of the nonexistence of $E$.

Theorem. [Zuo IMRN-2007] The prepotential $F$ is at most quadratic in the parameters $k_{\alpha}$, that is, up to quadratic terms in the flat coordinates

$$
F(\mathbf{t}, \mathbf{b})=F^{(0)}(\mathbf{t})+\sum_{\alpha=1}^{M} k_{\alpha} F^{(1)}\left(\mathbf{t}, b_{\alpha}\right)+\sum_{\alpha \neq \beta}^{M} k_{\alpha} k_{\beta} F^{(2)}\left(b_{\alpha}, b_{\beta}\right)
$$

where $\mathbf{t}=\left(t_{1}, \cdots, t_{l}\right)$ and $\mathbf{b}=\left(b_{1}, \cdots, b_{M}\right)$. Here $F^{(0)}$ is the potential associated to $B_{n}\left(D_{n}\right)$ and
$F^{(2)}\left(b_{\alpha}, b_{\beta}\right)=\frac{1}{2}\left(b_{\alpha}-b_{\beta}\right)^{2} \log \left(b_{\alpha}-b_{\beta}\right)^{2}+\frac{1}{2}\left(b_{\alpha}+b_{\beta}\right)^{2} \log \left(b_{\alpha}+b_{\beta}\right)^{2}$,
$\operatorname{deg} F=\operatorname{deg} F^{(0)}=4 K+2, \operatorname{deg} F^{(1)}=2 K+2, \operatorname{deg} F^{(2)}=2$.
$\S 4$. Frobenius manifolds and Extended affine Weyl groups

Motivation. Quantum cohomology of $\mathbb{P}^{1}$ :

$$
F=\frac{1}{2} t_{1}^{2} t_{2}+e^{t_{2}}, E=t_{1} \partial_{1}+2 \partial_{2}, e=\partial_{1}, W(M)=\widetilde{W}\left(A_{1}\right)
$$

Question: How to construct this kind of Frobenius manifolds? That is,

$$
\begin{aligned}
& F=F\left(t_{1}, \cdots, t_{n}, t_{n+1}, e^{t_{n+1}}\right) \\
& E=\sum_{\alpha=1}^{n} d_{\alpha} t_{\alpha} \partial_{\alpha}+d_{n+1} \partial_{n+1}
\end{aligned}
$$

## Notations

Let $R$ be an irreducible reduced root system defined on $(V,()$,$) .$ $\left\{\alpha_{j}\right\}$ : a basis of simple roots, $\quad\left\{\alpha_{j}^{\vee}\right\}$ : the corresponding coroots. $W$ Weyl group, $\quad W_{a}(R)$ affine Weyl group (the semi-direct product of $W$ by the lattice of coroots)
$W_{a}(R) \curvearrowright V$ : affine transformations

$$
\mathbf{x} \mapsto w(\mathbf{x})+\sum_{j=1}^{\prime} m_{j} \alpha_{j}^{\vee}, \quad w \in W, m_{j} \in \mathbb{Z}
$$

$\omega_{j}$ : the fundamental weights, $\left(\omega_{i}, \alpha_{j}^{\vee}\right)=\delta_{i j}$

Definition.[B.Dubrovin, Y.Zhang 1998]
The extended affine Weyl group $\widetilde{W}=\widetilde{W}^{(k)}(R)$ acts on the extended space

$$
\widetilde{V}=V \oplus \mathbb{R}
$$

and is generated by the transformations

$$
x=\left(\mathbf{x}, x_{I+1}\right) \mapsto\left(w(\mathbf{x})+\sum_{j=1}^{\prime} m_{j} \alpha_{j}^{\vee}, x_{I+1}\right), \quad w \in W, m_{j} \in \mathbb{Z}
$$

and

$$
x=\left(\mathbf{x}, x_{l+1}\right) \mapsto\left(\mathbf{x}+\gamma \omega_{k}, x_{l+1}-\gamma\right)
$$

Here $\gamma=1$ except for the cases when $R=B_{l}, k=I$ and $R=F_{4}, k=3$ or $k=4$, in these three cases $\gamma=2$.

Definition.[B.Dubrovin, Y.Zhang 1998]
$\mathcal{A}=\mathcal{A}^{(k)}(R)$ is the ring of all $\widetilde{W}$-invariant Fourier polynomials of the form

$$
\sum_{m_{1}, \ldots, m_{l+1} \in \mathbb{Z}} a_{m_{1}, \ldots, m_{l+1}} e^{2 \pi i\left(m_{1} x_{1}+\cdots+m_{l} x_{l}+\frac{1}{f} m_{l+1} x_{l+1}\right)}
$$

that are bounded in the limit

$$
\mathbf{x}=\mathbf{x}^{0}-i \omega_{k} \tau, \quad x_{l+1}=x_{l+1}^{0}+i \tau, \quad \tau \rightarrow+\infty
$$

for any $x^{0}=\left(\mathbf{x}^{0}, x_{l+1}^{0}\right)$, where $f$ is the determinant of the Cartan matrix of the root system $R$.

We introduce a set of numbers

$$
d_{j}=\left(\omega_{j}, \omega_{k}\right), \quad j=1, \ldots, l
$$

and define the following Fourier polynomials

$$
\begin{gathered}
\tilde{y}_{j}(x)=e^{2 \pi i d_{j} x_{l+1}} y_{j}(\mathbf{x}), \quad j=1, \ldots, I \\
\tilde{y}_{l+1}(x)=e^{\frac{2 \pi i}{\gamma} x_{l+1}}
\end{gathered}
$$

Here

$$
\begin{gathered}
y_{j}(\mathbf{x})=\frac{1}{n_{j}} \sum_{w \in W} e^{2 \pi i\left(\omega_{j}, w(\mathbf{x})\right)}, \\
n_{j}=\#\left\{w \in W \mid e^{2 \pi i\left(\omega_{j}, w(\mathbf{x})\right)}=e^{2 \pi i\left(\omega_{j}, \mathbf{x}\right)}\right\} .
\end{gathered}
$$

B.Dubrovin and Y.Zhang considered a particular choice of $\alpha_{k}$ based on the following observations

1. The Dynkin graph of $R_{k}:=\left\{\alpha_{1}, \cdots, \hat{\alpha_{k}}, \cdot, \alpha_{l}\right\}\left(\alpha_{k}\right.$ is omitted) consists of 1,2 or 3 branches of $A_{r}$ type for some $r$.
2. $d_{k}>d_{s}, s \neq k$.

Chevalley-Type Theorem [B.Dubrovin, Y.Zhang 1998]
For the above particluar choice of $\alpha_{k}$,

$$
\mathcal{A}^{(k)}(R) \simeq \mathbb{C}\left[\tilde{y}_{1}, \cdots, \tilde{y}_{l+1}\right] .
$$


$B_{I}$

$C_{1}$

$D_{I}$


$\begin{array}{lllll}F_{4} & \bullet & 0 & \longleftrightarrow & 0 \\ \\ & 1 & 2 & 3 & 4 \\ G_{2} & & \rightleftarrows & * \\ & & 1 & 2 & 3\end{array}$
$\mathcal{M}=\operatorname{Spec} \mathcal{A}$ : the orbit space of $\widetilde{W}^{(k)}(R)$ global coordinates on $\mathcal{M}:\left\{\tilde{y}_{1}(x), \cdots, \tilde{y}_{I+1}(x)\right\}$ local coordinates on $\mathcal{M}$ :

$$
y^{1}=\tilde{y}_{1}, \ldots, y^{\prime}=\tilde{y}_{l}, y^{l+1}=\log \tilde{y}_{l+1}=2 \pi i x_{l+1}
$$

the metric $(,)^{\sim}$ on $\widetilde{V}=V \oplus \mathbb{C}$
$\left(d x_{a}, d x_{b}\right)^{\sim}=\frac{1}{4 \pi^{2}}\left(\omega_{a}, \omega_{b}\right)$,
$\left(d x_{I+1}, d x_{a}\right)^{\sim}=0, \quad 1 \leq a, b \leq I$,
$\left(d x_{I+1}, d x_{I+1}\right)^{\sim}=-\frac{1}{4 \pi^{2}\left(\omega_{k}, \omega_{k}\right)}=-\frac{1}{4 \pi^{2} d_{k}}$
$\rightsquigarrow\left(\mathcal{M} \backslash \Sigma, g^{i j}(y)\right)$,

$$
\begin{equation*}
g^{i j}(y):=\left(d y^{i}, d y^{j}\right)^{\sim}=\sum_{a, b=1}^{l+1} \frac{\partial y^{i}}{\partial x^{a}} \frac{\partial y^{j}}{\partial x^{b}}\left(d x^{a}, d x^{b}\right)^{\sim} . \tag{0.2}
\end{equation*}
$$

Claim: $g^{i j}(y)$ is flat. Moreover for the particular choice, $g^{i j}(y)$ are at most linear w.r.t $y^{k}$.
$\rightsquigarrow \eta^{i j}(y)=\mathcal{L}_{e} g^{i j}(y)=\frac{\partial g^{i j}(y)}{\partial y^{k}}, \quad e:=\frac{\partial}{\partial y^{k}}$.

Theorem. [B.Dubrovin, Y.Zhang 1998]
For the particular choice of $\alpha_{k}, \eta^{i j}(y)$ and $g^{i j}(y)$ form a flat pencil. Moreover there exists a unique Frobenius structure on the orbit space $\mathcal{M}=\mathcal{M}(R, k)$ polynomial in $t^{1}, \ldots, t^{\prime}, e^{t^{\prime+1}}$ such that

1. the unity vector field coincides with $\frac{\partial}{\partial y^{k}}=\frac{\partial}{\partial t^{k}}$;
2. the Euler vector field has the form

$$
E=\frac{1}{2 \pi i d_{k}} \frac{\partial}{\partial x_{l+1}}=\sum_{\alpha=1}^{l} \frac{d_{\alpha}}{d_{k}} t^{\alpha} \frac{\partial}{\partial t^{\alpha}}+\frac{1}{d_{k}} \frac{\partial}{\partial t^{I+1}}
$$

3. The intersection form of the Frobenius structure coincides with the metric $(,)^{\sim}$ on $\mathcal{M}$.

Theorem.[P.Slodowy 1998,Preprint but unpublished]
The ring $\mathcal{A}^{(k)}(R)$ is isomorphic to the ring of polynomials of $\tilde{y}_{1}(x), \cdots, \tilde{y}_{l+1}(x)$ for arbitrary choice.

Another proof [B.Dubrovin, Y.Zhang and D.Zuo 2006]
We give an alternative proof of Chevelly-Type theorem associated to the root system $B_{l}, C_{l}, D_{l},\left(F_{4}, G_{2}\right)$.

Our question and result-II

An natural question:[P.Slodowy, B.Dubrovin and Y.Zhang 1998]
Is whether the geometric structures that were revealed in the above for particular choice also exist on the orbit spaces of the extended affine Weyl groups for an arbitrary choice of $\alpha_{k}$ ?

Difficulty: $d_{k}$ will be not the maximal number except the particular choice.

1. Note that the $g^{i j}(y)$ may be not linear with respect to $y^{k}$. Thus if we define $\eta^{i j}(y)=\frac{\partial g^{i j}(y)}{\partial y^{k}}$ as before, we can not obtain the flat pencil.
2. If we can obtain a flat pencil, how to find flat coordinates and construct Frobenius manifolds?

For the question 1 , our strategy is to change the unity vector field.
Main theorem 1. For any fixed integer $0 \leq m \leq I-k$ there is a flat pencil of metrics $\left(g^{i j}(y)\right),\left(\eta^{i j}(y)\right)$ (bilinear forms on $T^{*} M$ ) with $\left(g^{i j}(y)\right)$ given by (??) and $\eta^{i j}(y)=\mathcal{L}_{e} g^{i j}(y)$ on the orbit space $\mathcal{M}$ of $\widetilde{W}^{(k)}\left(C_{l}\right)$. Here the unity vector field

$$
e:=\sum_{j=k}^{l} a_{j} \frac{\partial}{\partial y^{j}}
$$

is defined by the generating function

$$
\sum_{j=k}^{I} a_{j} u^{I-j}=(u+2)^{m}(u-2)^{I-k-m}
$$

for the constants $a_{k}, \ldots, a_{l}$.

For the question 2, it is very technical.
Main theorem 2. In the flat coordinates $t^{1}, \ldots, t^{\prime+1}$, the nonzero entries of the matrix ( $\eta^{i j}$ ) are given by

$$
\eta^{i j}=\left\{\begin{array}{lll}
k, & j=k-i, & 1 \leq i \leq k-1 \\
1, & i=I+1, j=k & \text { or } i=k, j=I+1 \\
C, & j=I-m+k-i+1, & k+2 \leq i \leq I-m-1 \\
2, & i=I-m, j=k+1 & \text { or } i=k+1, j=I-m, \\
4 m, & j=2 I-m-i+1, & I-m+2 \leq i \leq I-1 \\
2, & i=I, j=I-m+1 & \text { or } i=I-m+1, j=I
\end{array}\right.
$$

where $C=4(I-m-k)$. The entries of the matrix $\left(g^{i j}(t)\right)$ and the Christoffel symbols $\Gamma_{m}^{i j}(t)$ are weighted homogeneous polynomials in $t^{1}, \ldots, t^{\prime}, \frac{1}{t^{\prime-m}}, \frac{1}{t^{\prime}}, e^{t^{\prime+1}}$.

Main theorem 3. For any fixed integer $0 \leq m \leq I-k$, there exists a unique Frobenius structure of charge $d=1$ on the orbit space $\mathcal{M} \backslash\left\{t^{I-m}=0\right\} \cup\left\{t^{\prime}=0\right\}$ of $\widetilde{W}^{(k)}\left(C_{l}\right)$ weighted homogeneous polynomial in $t^{1}, t^{2}, \cdots, t^{\prime}, \frac{1}{t^{\prime-m}}, \frac{1}{t^{\prime}}, e^{t^{\prime+1}}$ such that

1. The unity vector field e coincides with $\sum_{j=k}^{l} a_{j} \frac{\partial}{\partial y^{j}}=\frac{\partial}{\partial t^{k}}$;
2. The Euler vector field has the form

$$
E=\sum_{\alpha=1}^{l} \tilde{d}_{\alpha} t^{\alpha} \frac{\partial}{\partial t^{\alpha}}+\frac{\partial}{\partial t^{\prime+1}}
$$

3. The intersection form of the Frobenius structure coincides with the metric $\left(g^{i j}(t)\right)$.

Main theorem 4. The Frobenius manifold structures that we obtain in this way from $B_{I}$ and $D_{l}$, by fixing the $k$-th vertex of the corresponding Dynkin diagram, are isomorphic to the one that we obtain from $C_{l}$ by choosing the $k$-th vertex of the Dynkin diagram of $C_{1}$.

Example. [ $C_{5}, k=1, m=2$ ]Let $R$ be the root system of type $C_{5}$, take $k=1, m=2$, then

$$
\begin{aligned}
F= & \frac{1}{2} t_{6} t_{1}^{2}+\frac{1}{2} t_{1} t_{2} t_{3}+\frac{1}{2} t_{1} t_{4} t_{5}-\frac{1}{72} t_{3}{ }^{4} t_{5}^{4}-\frac{1}{8} t_{2} t_{3} t_{4} t_{5} \\
& -\frac{1}{2268} t_{5}^{8}-\frac{1}{36288} t_{3}{ }^{8}-\frac{1}{48} t_{3}{ }^{2} t_{2}{ }^{2}-\frac{1}{48} t_{4}{ }^{2} t_{5}{ }^{2}+\frac{1}{24} t_{5}{ }^{4} t_{2} t_{3} \\
& +\frac{1}{96} t_{3}{ }^{4} t_{4} t_{5}+\frac{1}{1440} t_{3}^{5} t_{2}+\frac{1}{360} t_{4} t_{5}{ }^{5}+t_{2} t_{3} e^{t_{6}}-t_{4} t_{5} e^{t_{6}} \\
& -\frac{2}{3} t_{5}{ }^{4} e^{t_{6}}+\frac{1}{6} t_{3}{ }^{4} e^{t_{6}}+\frac{1}{2} e^{2 t_{6}}+\frac{1}{48} \frac{t_{2}{ }^{3}}{t_{3}}+\frac{1}{192} \frac{t_{4}{ }^{3}}{t_{5}} .
\end{aligned}
$$

The Euler vector field is given by

$$
E=t_{1} \partial_{1}+\frac{3}{4} t_{2} \partial_{2}+\frac{1}{4} t_{3} \partial_{3}+\frac{3}{4} t_{4} \partial_{4}+\frac{1}{4} t_{5} \partial_{5}+\partial_{6} .
$$

## §5. Recent developments

Theorem.[2007] For any fixed integer $1 \leq k<I$, there exists a unique Frobenius structure of charge $d=1$ on the orbit space $\mathcal{M}^{k, 2}$ of $\widetilde{W}^{k, 2}\left(A_{l}\right)$ such that the potential $F(t)=\tilde{F}(t)+\frac{1}{2}\left(t^{k+1}\right)^{2} \log \left(t^{k+1}\right)$, where $\tilde{F}(t)$ is a weighted homogeneous polynomial in $t^{1}, t^{2}, \cdots, t^{\prime}, e^{t^{\prime+1}}, e^{t^{\prime+2}-t^{\prime+1}}$, satisfying

1. The unity vector field e coincides with $\frac{\partial}{\partial y^{k+1}}+e^{k y^{\prime+1}} \frac{\partial}{\partial y^{k}}=\frac{\partial}{\partial t^{k}}$;
2. The Euler vector field has the form

$$
E=\sum_{\alpha=1}^{I} \tilde{d}_{\alpha} t^{\alpha} \frac{\partial}{\partial t^{\alpha}}+\frac{1}{k} \frac{\partial}{\partial t^{\prime+1}}+\frac{I}{k(I-k)} \frac{\partial}{\partial t^{\prime+2}} .
$$

$A_{1}$


## Thanks

Appendix. Main techniques to obtain flat coordinates
The first step: $y \rightarrow \tau$

$$
\begin{aligned}
\sum_{j=0}^{I} \theta^{j} u^{I-j}= & \sum_{j=0}^{I-m} \varpi^{j}(u+2)^{m}(u-2)^{I-m-j} \\
& -\sum_{j=I-m+1}^{l} \varpi^{j}(u+2)^{I-j}(u-2)^{j-k-1}
\end{aligned}
$$

where

The second step: $\tau \rightarrow z$

$$
\begin{aligned}
& z^{I+1}=\tau^{I+1}, z^{j}=\tau^{j}+p_{j}\left(\tau^{1}, \ldots, \tau^{j-1}, e^{\tau^{\prime+1}}\right), 1 \leq j \leq k \\
& z^{j}=\tau^{j}+\sum_{s=j+1}^{I-m} c_{s}^{j} \tau^{s}, \quad k+1 \leq j \leq I-k-m \\
& z^{j}=\tau^{j}+\sum_{s=j+1}^{I} h_{s}^{j} \tau^{s}, \quad I-k-m+1 \leq j \leq I
\end{aligned}
$$

where $p_{j}$ are some weighted homegeoneous polynomials and $c_{s}^{j}$ and $h_{s}^{j}$ are determined by the following function respectively

$$
\cosh \left(\frac{\sqrt{t}}{2}\right)\left(\frac{2 \sinh \left(\frac{\sqrt{t}}{2}\right)}{\sqrt{t}}\right)^{2 i-1}, \quad\left(\frac{\tanh (\sqrt{t})}{\sqrt{t}}\right)^{2 i-1} .
$$

The third step: $z \rightarrow w$

$$
\begin{aligned}
& w^{i}=z^{i}, \quad i=1, \ldots, k, I+1 \\
& w^{k+1}=z^{k+1}\left(z^{I-m}\right)^{-\frac{1}{2(I-m-k)}} \\
& w^{s}=z^{s}\left(z^{I-m}\right)^{-\frac{s-k}{I-m-k}}, s=k+2, \cdots, I-m-1 \\
& w^{I-m}=\left(z^{I-m}\right)^{\frac{1}{2(I-m-k)}} \\
& w^{I-m+1}=z^{I-m+1}\left(z^{I}\right)^{-\frac{1}{2 m}} \\
& w^{r}=z^{r}\left(z^{l}\right)^{-\frac{r+m-I}{m}}, r=I-m+2, \cdots, I-1 \\
& w^{I}=\left(z^{I}\right)^{\frac{1}{2 m}}
\end{aligned}
$$

The last step: $w \rightarrow t$

$$
\begin{aligned}
& t^{1}=w^{1}, \ldots, t^{k}=w^{k}, t^{I+1}=w^{I+1} \\
& t^{k+1}=w^{k+1}+w^{I-m} h_{k+1}\left(w^{k+2}, \ldots, w^{I-m-1}\right), \\
& t^{j}=w^{I-m}\left(w^{j}+h_{j}\left(w^{j+1}, \ldots, w^{I-m-1}\right)\right), k+2 \leq j \leq I-m-1, \\
& t^{I-m+1}=w^{I-m+1}+w^{\prime} h_{I-m+1}\left(w^{I-m+2}, \ldots, w^{I-1}\right), \\
& t^{s}=w^{I}\left(w^{s}+h_{s}\left(w^{s+1}, \ldots, w^{I-1}\right)\right), I-m+2 \leq s \leq I-1 \\
& t^{I-m}=w^{I-m}, \quad t^{\prime}=w^{I} .
\end{aligned}
$$

Here $h_{l-m-1}=h_{l-1}=0, h_{j}$ are weighted homogeneous polynomials of degree $\frac{k(I-m-j)}{I-m-k}$ for $j=k+1, \ldots, I-m-2$ and $h_{s}$ are weighted homogeneous polynomials of degree $\frac{k(I-s)}{m}$ for $s=I-m+2, \ldots, l-1$.

Due to the above construction, we can associate the following natural degrees to the flat coordinates

$$
\begin{aligned}
& \tilde{d}_{j}=\operatorname{deg} t^{j}:=\frac{j}{k}, \quad 1 \leq j \leq k, \\
& \tilde{d}_{s}=\operatorname{deg} t^{s}:=\frac{2 l-2 m-2 s+1}{2(I-m-k)}, \quad k+1 \leq s \leq I-m, \\
& \tilde{d}_{\alpha}=\operatorname{deg} t^{\alpha}:=\frac{2 l-2 \alpha+1}{2 m}, \quad I-m+1 \leq \alpha \leq I, \\
& \tilde{d}_{l+1}=\operatorname{deg} t^{I+1}:=0 .
\end{aligned}
$$

